

Reese Chs. 7 & 8 (pp. 301-302 and 335-352)

- Motion in a Smooth Tunnel through Earth
- Gravitational Potential Energy
- Work-Energy Theorem
- Escape Speeds and Black Holes
- Limitations of Work-Energy Theorem

Motion in a Smooth Tunnel: I

- We imagine drilling a straight and smooth-walled tunnel along a diameter of a (hypothetical) uniform-density Earth.
- Applying Newton's inverse-square gravitation and his 2nd Law of Motion, we find that a rock released from rest at one end of the tunnel executes simple harmonic motion about Earth's center!

$$T = 2\pi \left(\frac{R_{\oplus}^3}{GM_{\oplus}} \right)^{\frac{1}{2}}$$

Motion in a Smooth Tunnel: II

- This period is exactly the same as the orbital period of an Earth-skimming satellite – about 84 minutes.
- Now imagine a straight smooth tunnel between two points on Earth's surface that are not at opposite ends of a diameter. Analysis gives SHM w/ same period as before!
- Hence "through the Earth in 42 minutes"!

Motion in Smooth Tunnel: III

- What happens if we are (slightly) more realistic and allow for the fact that the density increases toward Earth's center?
- Tunnel through Earth's center:
 - Is period less than, equal to, or greater than 84 minutes?
 - Is motion still simple harmonic motion?
- What about off-center tunnels?

Work (Review)

- In PHYS 240, we defined work done by a force along a given path as

$$W_{\text{path}} \equiv \int_{\text{path}} F \cos \theta dl = \int_{\text{path}} \vec{F} \cdot d\vec{r}$$

- In general, the work done getting from point P_i to P_f depends on the path.

Conservative Forces (Review)

- For certain forces, the work done by the force as a particle moves from one point to another is independent of the path. These forces are called **conservative forces**. The work done by a conservative force is zero for any closed path: $\oint_{\text{any closed path}} \vec{F}_{\text{conservative}} \cdot d\vec{r} = 0$

Additional Results I (Review)

- Net Work- Kinetic Energy Theorem: The work done by the net force on a particle equals the change in the particle's kinetic energy.
- When several forces are acting, the net work equals the sum of the individual works.
- For a conservative force, we can define a function of position called the potential energy. We begin by choosing a reference position P_0 where the potential energy is zero.

$$PE_a(\vec{r}) \equiv - \int_{\vec{r}_0}^{\vec{r}} \vec{F}_a(\vec{r}') \cdot d\vec{r}'$$

Additional Results II (Review)

- Notice the minus sign in the definition of potential energy. This means that the change in potential energy as the particle goes from P_i to P_f is equal to MINUS the work done by the conservative force:

$$PE_a(\vec{r}_f) - PE_a(\vec{r}_i) = - \int_{\vec{r}_i}^{\vec{r}_f} \vec{F}_a(\vec{r}') \cdot d\vec{r}'$$

Additional Results III (Review)

- Putting this all together, we find that if a particle is subject to some conservative forces ($a, b, c, \dots$) and some non-conservative forces (e.g. friction), then the net-work kinetic energy theorem implies that

$$KE_f - KE_i = W_{cons} + W_{non-cons}$$

$$\Rightarrow KE_f + \sum_{a, b, c, \dots} PE(\vec{r}_f) = KE_i + \sum_{a, b, c, \dots} PE(\vec{r}_i) + W_{non-cons}$$

Gravitational Potential Energy

- For a mass M at the origin, the Newtonian gravitational force on a test mass m at location $\vec{r}$ is

$$\vec{F}(\vec{r}) = -\frac{GMm}{r^2} \hat{r}$$

- It's not hard to show that this is a conservative force, and if we choose the location of zero potential energy to be $r = r_0$, then the gravitational potential energy function is

$$PE(\vec{r}) = -GMm \left(\frac{1}{r} - \frac{1}{r_0} \right)$$

Gravitational Potential of a System consisting of Earth and an Object near Earth's Surface: I

- For this system, the potential energy function is:

$$PE(\vec{r}) = -GM_{\oplus}m \left(\frac{1}{r} - \frac{1}{r_0} \right)$$

- If we choose the reference location P_0 to be the Earth's surface and we write r in terms of the object's altitude h , we can write the potential energy of the system as:

$$PE(h) = -GM_{\oplus}m \left[\left(\frac{1}{R_{\oplus} + h} \right) - \left(\frac{1}{R_{\oplus}} \right) \right]$$

Gravitational Potential of a System consisting of Earth and an Object near Earth's Surface: II

- Escape?: Notice for later use that the change in potential energy going from $h = 0$ to $h \rightarrow \infty$ is

$$\Delta(PE) = +\frac{GM_{\oplus}m}{R_{\oplus}}$$

- Low-altitude approximation: However, if $h \ll R_{\oplus}$ (radius of Earth), then

$$PE(h) = -\left(\frac{GM_{\oplus}m}{R_{\oplus}} \right) \left[\left(1 + \frac{h}{R_{\oplus}} \right)^{-1} - 1 \right] \approx \frac{GM_{\oplus}m}{R_{\oplus}^2} h = mg_{\oplus} h$$

Escape from Earth: I

- Imagine launching a projectile vertically from an idealized Earth (airless, nonrotating, spherical, isolated). What minimum initial speed gives escape?
- Apply the conservation of mechanical energy, taking final radius as infinite and final KE as 0:

$$\frac{1}{2}mv_i^2 + PE_\infty = \frac{1}{2}mv_i^2 + PE_i \Rightarrow 0 + (PE_\infty - PE_i) = \frac{1}{2}mv_{\text{escape}}^2$$

$$\Rightarrow v_{\text{escape}} = \sqrt{\frac{2GM_\oplus}{R_\oplus}}$$

Escape from Earth: II

- How does the minimum speed for escape depend on the launch direction?
 - It doesn't! Although the trajectory would be different, the final outcome is still escape.
- Suppose that an object is already in an Earth-skimming orbit. How much must its speed be increased to cause escape?

$$\Delta v_{\text{required}} = v_{\text{escape}} - v_{\text{loworbit}} = (\sqrt{2} - 1)v_{\text{loworbit}}$$

- The satellite's speed needs a 41.4% boost.

Escape from Anywhere? Nope!

- According to Newton, we can escape from the surface of Planet X as long as our launch speed exceeds

$$v_{\text{esc}X} = \sqrt{\frac{2GM_X}{R_X}}$$

- But Newton's gravity is inconsistent with Einstein's 1905 special relativity. Einstein's 1915 replacement gravity theory (general relativity) implies that nothing (not even light!) can escape from an object of mass M that has a radius less than

$$R_S \equiv \frac{2GM}{c^2}$$

Limitations on Use of Work-Energy Theorem

- For systems consisting of one or several simply pointlike particles for which there are no internal motions, deformations, or thermal effects, the work-energy theorem is a very useful tool.
- In systems with internal motions, deformations, and thermal effects, the work-energy theorem doesn't hold, although Newton's 2nd Law still applies to each part.
- Physicists do believe in and use conservation of total energy, but there are energy exchanges not covered by our macroscopic mechanics of rigid particles. See Reese pp. 351-352; for energy conservation, see Reese Ch.13.