

Review: Electrostatic Force

and Electrostatic Potential Energy

Review: Electrostatic Field

and Electrostatic Potential

Electric Potential due to

Continuously Distributed Charges

- · Equipotential Surfaces and Volumes
- · Field as the Negative Gradient of the Potential
- Motion of a Charge in E Field and the electron-volt

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Review: Electrostatic Force and Electrostatic Potential Energy

 The force on a test charge due to a static distribution of other electric charges is a <u>conservative</u> force, so that it is possible to define a corresponding electrostatic potential energy:

 $PE(\vec{r}_f) - PE(\vec{r}_i) \equiv -\int_{0}^{f} \vec{F}_{elec} \cdot d\vec{r}$

• NOTES:

- (1) Note the minus sign in the definition.
- (2) The equation above only determines <u>changes</u> in potential energy. The reference location (at which PE = 0) is a matter of choice.

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Review: Electrostatic Field and Electrostatic Potential

- The force on a test particle equals test particle's charge TIMES the local electrostatic field: \vec{F}
- Hence we can define an <u>electrostatic potential</u> V
 by saying that the change in electrostatic PE is
 just the test particle's charge times the change in
 the electrostatic potential. This implies that:

$$V(\vec{r}_f) - V(\vec{r}_i) \equiv -\int_{-1}^{f} \vec{E} \cdot d\vec{r}$$

Note minus sign and free choice of reference.

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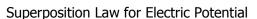
 The SI unit for electric potential is called the volt. It is equal to the SI unit of energy divided by the SI unit of charge:

1 $volt = 1 joule / coulomb \implies 1 V \equiv 1 J/C$

 NOTE: Be aware that the abbreviation for this unit (a capital V) is also commonly used as the variable name for electric potential!

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- The law of (vector) superposition for electric fields implies a (scalar) superposition law for electric potential.
- NOTE: In order to easily determine the reference (V=0) position when potentials are added, it's good practice to use the same reference position for all terms in the sum. For example, if we have a set of several fixed charges and adopt the choice that V=0 for r→∞, then

 $V(\vec{r}) = \frac{1}{4\pi\varepsilon_o} \sum_{n} \frac{Q_n}{|\vec{r} - \vec{r_n}|}$

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Electric Potential due to Continuously Distributed Charges

• When the charges are distributed continuously, the superposition law takes the form of an integral:

$$V(\vec{r}) = \frac{1}{4\pi\varepsilon_o} \int \frac{dQ}{\left|\vec{r} - \vec{r}_{dQ}\right|}$$

 RESTRICTION: The above equation can be used only when the charge distribution is confined to a bounded region. With infinitely extended distributions, we must calculate the electric field and then integrate the field to find potential differences.

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Equipotential Surfaces and Volumes

 From the definition of the electric potential in terms of the electric field, we find:

$$V(\vec{r} + d\vec{r}) - V(\vec{r}) \equiv dV \equiv -\vec{E}(\vec{r}) \cdot d\vec{r}$$

- This means that starting from any position vector ${\bf r}$, the neighboring points ${\bf r}$ + ${\bf dr}$ that have the same potential are those for which $\vec E \cdot {\bf d} \vec r = 0$:
 - (1) If the local field is nonzero, this defines an <u>equipotential surface</u>.
 - (2) If the local field is zero, then there can be a 3-D <u>equipotential volume</u>.

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Field as Negative Gradient of Potential: ${\bf I}$

• For <u>any</u> function of position, the <u>directional derivative</u> is defined by

$$\frac{df(\vec{r}, \hat{n})}{ds} = \lim_{\Delta s \to 0} \left[\frac{f(\vec{r} + \hat{n}\Delta s) - f(\vec{r})}{\Delta s} \right]$$

• It is tedious but not difficult to show that:

$$\boxed{\frac{df(\vec{r},\hat{n})}{ds} = \hat{n} \cdot \left[\hat{i} \frac{\partial f}{\partial x} + \hat{j} \frac{\partial f}{\partial y} + \hat{k} \frac{\partial f}{\partial z} \right] \equiv \hat{n} \cdot \nabla f}$$

 $ec{
abla}f$ is called the gradient of the function f.

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- Examining how the electric field is related to the potential function, and
- comparing that with how the gradient of f is related to f, we conclude that the electric field is nothing but the negative gradient of the electric potential:

$$\vec{E}(\vec{r}) = -\vec{\nabla}V(\vec{r})$$

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Chapter 1

Acceleration of Charges by Electrostatic Fields

• Newton's 2nd Law provides:

$$m\vec{a} = \vec{F}_{net} = q\vec{E} = -q\vec{\nabla}V$$

• It is easy to use this to derive:

$$\frac{1}{2}mv_f^2 + qV_f = \frac{1}{2}mv_i^2 + qV_i$$

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An auxiliary energy unit: the electron-volt

- An electron-volt is the kinetic energy increase exhibited by an electron when its electric potential increases by one volt. (Notice that since the electron has a negative charge, it is accelerated in the direction of increasing electric potential – opposite the direction in which the field points.)
- Since the magnitude of the charge on an electron is $1.602 \times 10^{-19} \, C$, we have

$$1 \ electron - volt = 1 \ eV = 1.602 \times 10^{-19} J$$

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Reese Ch. 17, part B (pp. 786-795)

- Review: An Electric Dipole & its Field (pp. 728-733)
- An Electric Dipole in an External Field
- Electric Potential due to a Dipole
- Using Potential to find the Field of a Dipole
- Potential Energy of an Assembly of Point Charges
- · Lightning Rods

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- An <u>electric dipole</u> consists of two charges of equal magnitude |O| but opposite sign separated by a relatively short distance d.
- The dipole moment **p** of a dipole is a vector with magnitude |Q| times d, located midway between the charges, and pointing from the negative charge toward the positive one.
- For r >> d, $|\mathbf{E}|$ varies inversely as r^3 .

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An Electric Dipole in an External Field

- If an electric dipole is placed in a uniform external electric field **E**, there is zero net force on the dipole (Why?).
- There is a net torque on the dipole: (Note that because the net force is zero, this expression for the torque is correct for any choice of coordinate origin.)
- In a non-uniform electric field, a dipole suffers both a torque and a net force. (Why?)

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Electric Potential due to a Dipole

• At radii r >> d, the electric potential due to a dipole $\vec{p} = (|Q|d)k$ centered at the origin is given to a very good approximation by:

$$V_{dipole}(r,\theta) = \frac{|Q|}{4\pi\varepsilon_o} \frac{d\cos\theta}{r^2} = \frac{1}{4\pi\varepsilon_o} \frac{p\cos\theta}{r^2}$$

- - (1) Here the reference is V → 0 as r → ∞.
 - (2) θ is the colatitude (angle between **r** and +z)
 - (3) For given r, how does V vary with θ?

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Using the Potential to Find the Field

$$\vec{E}(\vec{r}) = -\vec{\nabla}V(\vec{r})$$

In spherical polar coordinates,

$$\vec{\nabla} f \equiv \hat{r} \frac{\partial f}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial f}{\partial \theta} + \hat{\varphi} \frac{1}{r \sin \theta} \frac{\partial f}{\partial \varphi}$$

$$\Rightarrow \vec{E}(r,\theta) = \frac{p}{4\pi\varepsilon_o} \left[\hat{r} \left(\frac{2\cos\theta}{r^3} \right) + \hat{\theta} \left(\frac{\sin\theta}{r^3} \right) \right]$$

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Potential Energy of an Assembly of Point Charges

 Using Coulomb's law, the work required to bring two charges near each other (from infinite separation) is:

$$W_{to}_{assemble} = \Delta(PE) = \frac{1}{4\pi\varepsilon_o} \frac{Q_1 Q_2}{|\vec{r_2} - \vec{r_1}|} \equiv \frac{1}{4\pi\varepsilon_o} \frac{Q_1 Q_2}{r_{12}}$$

 With infinite mutual separation as the zero for PE, we can generalize to N charged particles:

$$PE = \frac{1}{4\pi\varepsilon_o} \sum_{i=1}^{N} \sum_{j=i+1}^{N} \frac{Q_i Q_j}{|\vec{r}_j - \vec{r}_i|} = \frac{1}{4\pi\varepsilon_o} \sum_{i=1}^{N} \sum_{j=i+1}^{N} \frac{Q_i Q_j}{r_{ij}}$$

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- These are among the many useful inventions of Benjamin Franklin.
- They work because:
 - (1) The magnitude of the electric field immediately surrounding a conducting object is greatest where the object is "sharpest" (has smallest radius of curvature).
 - (2) Air breaks down (ionizes) and provides a conducting path when the field exceeds a certain critical value (about 3 x 10⁶ N/C = 3 x 10⁶ V/m).

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